

# High-Temperature Expansions to Fifteenth Order

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High-temperature series expansions of the susceptibility and second moment to 15th order are calculated for zero external field on the linear chain (LC), plane square (PSQ), simple cubic (SC), and body-centered cubic (BCC) lattices. Checks for specific models against pertinent work in the literature are detailed.

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**KEY WORDS:** High-temperature series; high-temperature expansions; strong coupling series; strong coupling expansions.

## 1. INTRODUCTION

High-temperature series expansions have proven to be a powerful method of studying lattice models. In particular, much attention has been devoted to the comparison of predictions for critical exponents derived from series expansions and from renormalization group methods (see, e.g., Refs. 1).

The unrenormalized lattice strong coupling expansion in quantum field theory is directly related to the high-temperature series expansion of statistical mechanics, with the square of the bare coupling playing the role of temperature. The expansions are lattice- and potential-dependent, of course, but the dependence on potential can be absorbed into vertex functions, which then can be left general, so that the remaining coefficients are only a function of lattice geometry.

In this paper, we use a diagrammatic prescription very similar to that of Bender *et al.*<sup>(2)</sup> to obtain the high-temperature (strong coupling) expansion of any scalar spin theory (quantum field theory) with nearest neighbor interactions with an even spin density (potential). We compute the susceptibility  $\chi$  and the second moment  $\mu_2$  in zero external field to 15th order in the expansion parameter for four different lattices: the linear chain (LC), the

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plane square lattice (PSQ), the simple cubic lattice (SC), and the body-centered cubic lattice (BCC).

The power of this method is derived from three characteristics of the approach:

1. The graph counting and combinatorial weighting have been completely automated, so that computer resources (and not human stamina and fallibility) limit the order of perturbative calculations. The graph counting method is new.
2. The absence of tadpoles (lines beginning and ending at the same vertex) greatly reduces the number of graphs that must be considered to a given order.
3. The spin-density-dependent (potential-dependent) information is entirely contained in the vertex functions, so that once the pertinent graphs have been enumerated, the series is applicable to any spin-density (potential) that is even.

In Section 2 we give a brief description of the expansion and how it may be reexpressed as a series of terms each of which factors into separate lattice-dependent and spin-density-dependent terms. In Section 3 we outline how one-vertex irreducible subgraphs (OVIS) may be used to generate the set of the one-particle irreducible graphs that are then assembled to form all graphs that contribute to the 15th-order expansion. The algebraic procedure used to form the susceptibility and second moment from the one-particle irreducible graphs is described in Section 4. The results, which extend the tenth-order tables of Kincaid *et al.*,<sup>(3)</sup> appear in a similar tabular form in Section 5. A discussion of the consistency checks made on them can be found in Section 6. The Appendix describes in some detail the new algorithm used to generate the OVIS.

The results of this paper have already been used to study hyperscaling in the Ising model.<sup>(4)</sup>

## 2. THE DIAGRAMMATIC EXPANSION

Using the language of quantum field theory, we begin with a generating functional (partition function) on a  $d$ -dimensional Euclidean lattice

$$Z[J] = \int [D\phi] \exp \sum_i a^d \left( \frac{\phi_i(\phi_{i+1} + \phi_{i-1} - 2d\phi_i)}{2a^2} - V(\phi_i) - J_i \phi_i \right) \quad (2.1)$$

which is the lattice version of the continuous generating functional

$$Z[J] = \int [D\phi] \exp \left\{ - \int d^4x \left[ \frac{1}{2} (\partial\phi)^2 + V(\phi) + J\phi \right] \right\} \quad (2.2)$$

The potential  $V(\phi_i)$  is a local one, which depends only on the value of the field at site  $i$ ;  $J_i\phi_i$  is a source term. The “off-diagonal” piece of the kinetic energy operator may be extracted from the action to act as an operator on the remaining portion of the functional integral,

$$\begin{aligned} Z[J] = & \exp \sum_i \frac{v}{2} \left[ \frac{\delta}{\delta J_i} \left( \frac{\delta}{\delta J_{i+1}} + \frac{\delta}{\delta J_{i-1}} \right) \right] \\ & \times \int [D\phi] \exp \left\{ - \sum_i a^d \left[ V(\phi_i) + \frac{d\phi_i^2}{a^2} + J_i\phi_i \right] \right\} \end{aligned} \quad (2.3)$$

with  $v = a^{-(d+2)}$ . This produces a formal expansion for the generating functional, from which the  $n$ -point Green's functions of the theory (correlation functions) can be obtained by a simple diagrammatic procedure.

In a strong coupling expansion, the diagrams are organized by increasing powers of  $v$ , the expansion parameter. This parameter counts the number of free inverse propagators (lines) in the graph, in contrast to the organization of graphs by increasing powers of  $g$  (vertices) in weak coupling expansions. The  $n$ -point connected Green's functions  $G_n(\phi_1, \phi_2, \dots, \phi_n)$  of the theory at zero source field can then be computed from  $Z(J)$  by taking functional derivatives with respect to the external source  $J(\phi)$ :

$$G_n(\phi_1, \phi_2, \dots, \phi_n) = \frac{\delta}{\delta J(\phi_1)} \cdots \frac{\delta}{\delta J(\phi_n)} \ln Z[J] \Big|_{J=0} \quad (2.4)$$

Normalization of the generating functional is of no concern, therefore, if we restrict our attention to Green's functions.

Since we consider only potentials that are even in the field variable  $\phi$ , all vertices represent the intersection of an even number of bare propagators. As in the usual high-temperature series expansion, we have only factored out the off-diagonal portion of the kinetic energy operator, and as a result there are no graphs containing tadpoles.

The  $n$ -point connected Greens' functions can be constructed from the formal expansion through a given order in  $v$  by the use of a set of diagrammatic rules. Following Bender *et al.*<sup>(2)</sup> we list these rules for constructing  $n$ -point Green's functions to order  $N$  in the internal line counter  $v$ :

1. Draw all connected graphs with no tadpoles having a total of  $n$  external lines and  $N$  internal lines. The vertices of these graphs must be the intersection of an even number of lines (external or internal). This enumeration of graphs is general to all even scalar theories in zero external field.

2. Associate with every  $2p$  vertex the amplitude  $\lambda_{2p}$  given by

$$\lambda_{2p} = v^p W_{2p}$$

where the  $W_{2p}$  are the vertex functions defined by

$$\exp \sum_{p=0}^{\infty} \frac{J^{2p} W_{2p}}{(2p)!} = \frac{\int d\phi \exp \{-a^d [V(\phi) + d\phi^2/a^2 + J\phi]\}}{\int d\phi \exp \{-a^d [V(\phi) + d\phi^2/a^2]\}} \quad (2.5)$$

The vertex functions carry all potential-dependent information, but are lattice-independent.

3. With each graph, associate two different coefficients: the topological symmetry number and the free multiplicity.

(a) The topological symmetry number is the inverse of the number of ways that an identical graph may be obtained by relabeling the internal lines and vertices of the graph. This number is potential- and lattice-independent.

(b) The free multiplicity counts the number of ways that a given graph with fixed external lines may be embedded onto the chosen lattice with the vertices corresponding to lattice sites and the internal lines connecting neighboring sites. This number is potential-independent but lattice-dependent and may be zero, in fact, for some graphs on some lattices.

4. Form the product of the vertex amplitude, the topological symmetry number, the free multiplicity, and  $v^{-n/2}$  to arrive at the total contribution to the series from each graph.

5. Sum the contributions of all graphs.

We implemented these rules through 15th order for several different lattice geometries on a DEC-10 using primarily FORTRAN, and to a lesser extent the symbolic manipulation routines ASHMEDEI<sup>(5)</sup> and REDUCE.<sup>(6)</sup> Before describing our technique for building all connected graphs from a smaller set of graphs, let us illustrate the rules above with an example.

Suppose we wish to compute the susceptibility [that is, the two-point correlation function  $G_2(x, y)$  when summed over all external points (y)] for an even scalar theory on a plane square lattice through third order in our perturbation parameter  $v$ . We begin by enumerating all graphs relevant to this order, i.e., those containing no more than three internal lines. They are shown in Fig. 1.

The topological symmetry number and free multiplicity are also shown for each of these graphs. Notice that the topological symmetry number depends only on the shape of the graph, but the free multiplicity depends

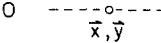
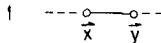
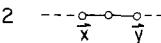
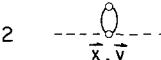
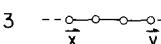
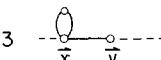
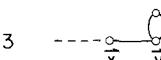
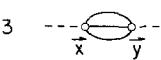
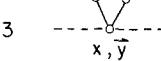
ORDER	GRAPH	TOPOLOGICAL SYMMETRY NUMBER	FREE MULTIPLICITY [plane square lattice]	$ \vec{x} - \vec{y} $	CONTRIBUTION TO $\chi \equiv \sum_{\vec{y}} G_2(\vec{x}, \vec{y})$
0		1	1	0	$w_2$
1		1	1	1	$4v w_2^2$
2		1	4 2 1	$\frac{0}{\sqrt{2}}$ $\frac{\sqrt{2}}{2}$	$16 v^2 w_2^3$
2		$\frac{1}{2}$	4	0	$2v^2 w_2 w_4$
3		1	9 3 1	$\frac{1}{\sqrt{5}}$ $\frac{\sqrt{5}}{3}$	$64 v^3 w_2^4$
3		$\frac{1}{2}$	4	1	$8v^3 w_2^2 w_4$
3		$\frac{1}{2}$	4	1	$8v^3 w_2^2 w_4$
3		$\frac{1}{6}$	1	1	$\frac{2}{3}v^3 w_4^2$
3		$\frac{1}{2}$	0	0	0

Fig. 1. All graphs contributing through third order to the strong coupling expansion of the two-point Green's function  $G_2(x, y)$ . External lines are dashed; internal lines are solid. The topological symmetry number of each graph is shown with the free multiplicity of the graph for different embeddings  $(x, y)$  on a plane square lattice. The total contribution of each graph to the susceptibility through third order is also given.

on the shape of the graph, the type of lattice, and the relative location of the external legs. For example, the graph with two internal lines and three vertices can be placed on a plane square lattice in only one way if the external points are  $(0, 0)$  and  $(2, 0)$ . But the same graph can be laid onto the plane square lattice in two different ways if the external points are  $(0, 0)$  and  $(1, 1)$ . Note, too, that the last graph in the list does not contribute on a plane square lattice, but would contribute on a plane triangular lattice.

Counting the number of internal lines to arrive at the appropriate

power of  $v$  is trivial, and the sum over one of the external points  $\mathbf{y}$  (up to three lattice units away from  $\mathbf{x}$ ) of all contributions then gives the susceptibility to third order.

### 3. ONE-VERTEX IRREDUCIBLE SUBGRAPHS

In order to automate the graph bookkeeping to the largest possible extent, we first build a subset of graphs, called the kernel. These graphs are constructed according to the rules of the previous section for the two-point function, with the added restriction that they be one-particle irreducible. One-particle irreducible graphs are those that cannot be separated into two disjoint graphs by the removal of a single internal line. From the kernel  $K(\mathbf{x}, \mathbf{y})$ , other physical quantities of interest, such as the full two-point function  $G_2(\mathbf{x}, \mathbf{y})$  may be built. The kernel itself is built from smaller building blocks, which we call one-vertex irreducible subgraphs (OVIS).<sup>(7)</sup> The set of OVIS is the set of all one-particle irreducible graphs in the two-point function for which the removal of any internal vertex will not separate the graph into two disjoint pieces only one of which contains both external legs. Such a set of graphs is, in fact, equivalent to the full kernel if the vertex functions are appropriately redefined. In this section we describe the algorithm for generating OVIS and the vertex redefinition.

All OVIS through sixth order in the internal line counter are shown in Fig. 2, with the redefined vertices indicated as dots centered in open circles. For each graph, the topological symmetry number is shown along with the free multiplicity for all possible embeddings on a plane square lattice. The degrees of a graph are the number of lines that enter each of its vertices. The sum of the degrees must equal  $2N + n$ , where  $N$  is the order of the graph and  $n$  is the number of external legs. In the Appendix we describe our new algorithm for generating all such OVIS to a given order together with their topological symmetry numbers and lattice-dependent free multiplicities.

Having generated the OVIS, the task of applying the appropriate vertex redefinition remains. The redefinition must build the full kernel by the addition of graphs that are one-vertex reducible. The redefined vertices, the  $M_{2p}$ , can then be expanded as a power series in the previously defined  $W_{2p}$ . This redefinition can be thought of pictorially as beginning with a “bare”  $W_{2p}$  vertex and attaching vacuum “decorations” that have an even number of internal lines and bare  $W_{2p}$  internal vertices. The sum of all possible decorations to a given order form the “decorated”  $M_{2p}$  vertex functions of the OVIS. Thus, the set of OVIS hide in their decorated vertices the elements of the kernel that are one-vertex reducible.

Shown in Fig. 3 is the pictorial expansion through fourth order of the

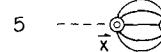
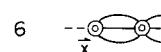
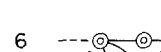
ORDER	GRAPH	TOPOLOGICAL SYMMETRY NUMBER	FREE MULTIPLICITY [plane square lattice]	$ \vec{x} - \vec{y} $	DEGREES
0		1	1	0	2
3		$\frac{1}{6}$	1	1	4,4
5		$\frac{1}{120}$	1	1	6,6
5		$\frac{1}{2}$	9	1	4,4,2,2
6		$\frac{1}{36}$	4 2 1	$\sqrt{2}$ $\sqrt{2}$ $\sqrt{2}$	6,4,4
6		$\frac{1}{4}$	16 4 1	$\sqrt{2}$ $\sqrt{2}$ $\sqrt{2}$	4,4,4,2
6		$\frac{1}{6}$	64 8 1	$\sqrt{2}$ $\sqrt{2}$ $\sqrt{2}$	4,4,2,2,2

Fig. 2. All one-vertex irreducible subgraphs contributing through sixth order to the strong coupling expansion. External lines are dashed; internal lines are solid. The redefined, or "decorated," vertices  $M_{2p}$  are shown as dots centered in open circles. The topological symmetry number of each graph is shown with the free multiplicity for different embeddings on a plane square lattice. The degrees are also shown for each graph.

decorated vertex of lowest degree,  $M_2$ , into its bare vertex decorations. The algebraic expansion is also given to the same order for comparison. Each decoration is of even order and is associated with its own topological symmetry number. The bare vertices of the decorations are all internal and their positions are therefore summed over freely to obtain the free multiplicity sum, which, although independent of the external positions  $x$  and  $y$ , is, of course, lattice-dependent. Notice that for decorations the sum of the degrees is again equal to  $2N+n$ , where  $N$  is now the number of internal lines added by the decoration and  $n$  is the degree of the decorated vertex.

We choose this redefinition for the vertex functions of the OVIS since a simple substitution of the series expansions in the bare  $W_{2p}$  for the decorated  $M_{2p}$  produces composite graphs for which the topological

DECORATION	ORDER	TOPOLOGICAL SYMMETRY NUMBER	FREE MULTIPLICITY [plane square lattice]	DEGREES
$\text{---} \circ \text{---} = \text{---} \circ \text{---}$ M2	0	1	1	2
$+ \frac{\text{---}}{2v^2} \frac{\text{---}}{W_4 W_2}$	2	$\frac{1}{2}$	4	4, 2
$+ \frac{\text{---}}{\frac{1}{6}v^4} \frac{\text{---}}{W_6 W_4}$	4	$\frac{1}{24}$	4	4, 4
$+ \frac{\text{---}}{2v^4} \frac{\text{---}}{W_6 W_2^2}$	4	$\frac{1}{8}$	16	6, 2, 2
$+ \frac{\text{---}}{4v^4} \frac{\text{---}}{W_4^2 W_2}$	4	$\frac{1}{4}$	16	4, 4, 2
$+ \frac{\text{---}}{18v^4} \frac{\text{---}}{W_4 W_2^3}$	4	$\frac{1}{2}$	36	4, 2, 2, 2

Fig. 3. All decorations contributing through fourth order to the decorated vertex  $M_2$ . Lines associated with the decorated vertex are dashed; lines internal to the decoration are solid. The decorated vertex is shown as a dot in an open circle; the bare vertices are shown as dots. The order of each decoration is shown with its topological symmetry number. The free multiplicity sum for each decoration is shown here for embeddings on a plane square lattice. The degrees for each decoration are also indicated.

symmetry numbers and free multiplicities are a simple product of those for the constituent OVIS and decorations. This procedure again lends itself to automation.

#### 4. THE SUSCEPTIBILITY $\chi$ AND THE SECOND MOMENT $\mu_2$

By expanding the decorated vertices of the set of OVIS to a given order, one generates the full one-particle irreducible kernel  $K(x, y)$ . From this kernel, the following then may be easily found: the two-point function

$G_2(\mathbf{x}, \mathbf{y})$ ; its sum over external  $\mathbf{y}$ , namely the susceptibility  $\chi = \sum_{\mathbf{y}} G_2(0, \mathbf{y})$ ; and the second moment of the two-point function  $\mu_2 = \sum_{\mathbf{y}} G_2(0, \mathbf{y}) \mathbf{y}^2$ .

The two-point function  $G_2(\mathbf{x}, \mathbf{y})$  is the sum of all graphs with vertices connected to external lines at  $(\mathbf{x}, \mathbf{y})$ . All such graphs can be formed from the kernel  $K(\mathbf{x}, \mathbf{y})$ , and the bare propagator  $P(\mathbf{x}, \mathbf{y})$  by forming the following series, which can then be truncated at the desired order:

$$G_2(\mathbf{x}, \mathbf{y}) = K(\mathbf{x}, \mathbf{y}) + \sum_{\mathbf{x}', \mathbf{x}''} K(\mathbf{x}, \mathbf{x}') P(\mathbf{x}', \mathbf{x}'') K(\mathbf{x}'', \mathbf{y}) + \dots \quad (4.1)$$

The bare propagator is simply related to the line counter  $v$ , since

$$P(\mathbf{x}, \mathbf{y}) = v \{ \delta_{\mathbf{x}, \mathbf{y} + \hat{\mathbf{u}}} + \delta_{\mathbf{x}, \mathbf{y} - \hat{\mathbf{u}}} \} \quad (4.2)$$

where  $\hat{\mathbf{u}}$  is a vector that lies along one of the positive lattice axes and connects neighboring lattice points.

The susceptibility  $\chi$  is the sum over one external point of the two-point function

$$\chi \equiv \sum_{\mathbf{y}} G_2(0, \mathbf{y}) \quad (4.3)$$

If we further define an integrated kernel  $K_0$  as

$$K_0 \equiv \sum_{\mathbf{y}} K(0, \mathbf{y}) \quad (4.4)$$

then, using the translation invariance of the lattice, we may write the susceptibility as the following power series in  $K_0$ :

$$\chi = \sum_{p=0}^{\infty} (2\eta v)^p K_0^{p+1} \quad (4.5)$$

where  $\eta$  is the number of axes of symmetry for the lattice. Since  $K_0$  itself is a power series in  $v$ , we must remember to keep in each term of Eq. (4.5) only those pieces that contribute to the desired order of expansion.

As an example, consider the susceptibility through third order in  $v$  on a plane, square lattice ( $\eta = 2$ ). The one-particle irreducible graphs contributing to the kernel at this order are shown in Fig. 4. The integrated kernel  $K_0$  is thus given by

$$K_0 = W_2 + 2v^2 W_2 W_4 + \frac{2}{3} v^3 W_4^2 \quad (4.6)$$

The reader may verify (against Fig. 1) that substitution of this expression into Eq. (4.5), keeping terms no higher than  $v^3$ , will reproduce the full set of graphs with their corresponding contribution to the susceptibility.

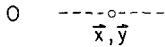
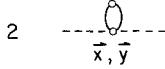
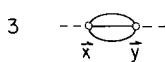
ORDER	GRAPH	TOPOLOGICAL SYMMETRY NUMBER	FREE MULTIPLICITY [plane square lattice]	$ x - \vec{y} $	CONTRIBUTION TO $K_0 \equiv \sum_{\vec{y}} K(0, \vec{y})$
0		1	1	0	$w_2$
2		$1_2$	4	0	$2v^2 w_2 w_4$
3		$1_6$	1	1	$\frac{2}{3} v^3 w_4^2$

Fig. 4. All kernel graphs through third order on a plane square lattice. External lines are dashed; internal lines are solid. The contribution of each graph to the integrated kernel  $K_0$  is shown.

The second moment  $\mu_2$  is defined as

$$\mu_2 \equiv \sum_y G_2(0, y) y^2 \quad (4.7)$$

and is not translationally invariant. It is related to the Fourier transform  $\tilde{G}_2$  of  $G_2$  by

$$\mu_2 = -\nabla_p^2 \tilde{G}_2(\mathbf{p})|_{\mathbf{p}=0} \quad (4.8)$$

From Eq. (4.1) we have

$$\tilde{G}_2(\mathbf{p}) = \tilde{K}(\mathbf{p}) + \tilde{K}(\mathbf{p}) \tilde{P}(\mathbf{p}) \tilde{K}(\mathbf{p}) + \dots \quad (4.9)$$

or

$$\tilde{G}_2(\mathbf{p}) = \frac{\tilde{K}(\mathbf{p})}{1 - \tilde{P}(\mathbf{p}) \tilde{K}(\mathbf{p})} \quad (4.10)$$

Now using the fact that the two-point function is an even function of  $\mathbf{p}$ , (lattice isotropy), we apply  $-\nabla_p^2$  to Eq. (4.10) to arrive at

$$\begin{aligned} \mu_2 &= \frac{-[\nabla_p^2 \tilde{K}(\mathbf{p})]}{1 - \tilde{P}(\mathbf{p}) \tilde{K}(\mathbf{p})} \Big|_{\mathbf{p}=0} \\ &= \tilde{K}(\mathbf{p}) \left\{ \frac{\tilde{P}(\mathbf{p}) [\nabla_p^2 \tilde{K}(\mathbf{p})] + [\nabla_p^2 \tilde{P}(\mathbf{p})] \tilde{K}(\mathbf{p})}{[1 - \tilde{P}(\mathbf{p}) \tilde{K}(\mathbf{p})]^2} \right\} \Big|_{\mathbf{p}=0} \end{aligned} \quad (4.11)$$

Note that

$$\begin{aligned}\tilde{P}(0) &= \sum_y P(0, y) = 2v\eta \\ \tilde{K}(0) &= \sum_y K(0, y) \equiv K_0 \\ -\nabla_p^2 \tilde{P}(0) &= \sum_y P(0, y) y^2 = 2v\eta\end{aligned}\quad (4.12)$$

and further, define  $K_2$  to be the second moment of the kernel, that is

$$-\nabla_p^2 \tilde{K}(\mathbf{p})|_{p=0} = \sum_y K(0, y) y^2 \equiv K_2 \quad (4.13)$$

So, as a power series in  $K_0$ ,

$$\mu_2 = \frac{K_2 + 2v\eta K_0^2}{(1 - 2v\eta K_0)^2} = (K_2 + 2v\eta K_0^2) \sum_{p=0}^{\infty} (p+1)(2v\eta K_0)^p \quad (4.14)$$

Since  $K_2$  begins at order  $v$ , so does the second moment  $\mu_2$ .

As an exercise, the reader may verify that through third order on a plane square lattice the second moment is

$$\mu_2 = 4vW_2^2 + 32v^2W_2^3 + v^3(192W_2^4 + 16W_2^2W_4 + \frac{2}{3}W_4^2) \quad (4.15)$$

Using ASHMEDAI and REDUCE, the series in Eqs. (4.5) and (4.14) were expanded algebraically and truncated at the 15th order for four different lattice geometries. The results are presented in tabular form at the end of the paper.

## 5. THE RESULTS

The results of our strong coupling expansion of the susceptibility  $\chi$  and second moment  $\mu_2$  are given in Tables I and II respectively. We have tabulated the series, order by order, in a form that is completely analogous to the tenth-order tables of Kincaid *et al.*<sup>(3)</sup> whose work we extend here through 15th order. Note that Kincaid *et al.* computed  $\chi$  and  $\mu_2$  on the more complicated triangular and face-centered cubic lattices, and also computed  $\partial^2\chi/\partial H^2$  for all lattice geometries considered; we did not.

The series may be reconstructed from the tables by multiplying each partition by the corresponding coefficient and summing through the desired order. The coefficients contain all the lattice-dependent information and are tabulated for the four lattice geometries of this work: linear chain

(LC), plane square (PSQ), simple cubic (SC), and body-centered cubic (BCC).

The partitions are grouped by leading power of the expansion parameter  $v$  and indicate (in ascending order) the degree of each strong coupling moment  $I_{2n}$  found in that term. The moments are related to the vertex functions  $W_{2n}$  through

$$\sum_{n=0}^{\infty} \frac{J^{2n} I_{2n}}{(2n)!} = \exp \sum_{n'=0}^{\infty} \frac{J^{2n'} W_{2n'}}{(2n')!} \quad (5.1)$$

and contain, therefore, all the potential-dependent information. By differentiating both sides of Eq. (5.1) with respect to  $J$ , we obtain

$$\sum_{n=1}^{\infty} \frac{I_{2n} J^{2n-1}}{(2n-1)!} = \left( \sum_{n'=1}^{\infty} \frac{W_{2n'} J^{2n'-1}}{(2n'-1)!} \right) \left( \sum_{m=0}^{\infty} \frac{I_{2m} J^{2m}}{(2m)!} \right) \quad (5.2)$$

Now, by collecting like terms in powers of  $J$ , we arrive at the following recursive formula for the expansion of the  $I_{2n}$  as a power series in the  $W_{2n}$ :

$$I_{2n} = (2n-1)! \left\{ \sum_{n'=1}^n \frac{W_{2n'}}{(2n'-1)!} \frac{I_{2(n-n')}}{[2(n-n')]!} \right\} \quad (5.3)$$

Note that for convenience, and to ensure that all coefficients listed in the tables are integers, we have multiplied each coefficient in the leading  $n$ th order by  $n!$ .

To illustrate the format via an example, consider the susceptibility on a plane square lattice to third order. Table I indicates the following series in terms of the moments  $I_{2n}$ :

$$\chi = I_2 + 4vI_2^2 + \frac{v^2}{2!} (20I_2^3 + 4I_2 I_4) + \frac{v^3}{3!} (132I_2^4 + 72I_2^2 I_4 + 4I_4^2) \quad (5.4)$$

Using Eq. (5.3), we find

$$I_2 = W_2, \quad I_4 = W_4 + 3W_2^2 \quad (5.5)$$

The reader may verify that substitution yields the expression for the susceptibility on a plane square lattice through third order previously found in Fig. 1.

The complete set of tables is available from the authors in computer-ready form, via electronic mail or magnetic tape.

## 6. VERIFICATION

Our series extend an existing body of literature in strong coupling expansions in field theory and high-temperature expansions in statistical mechanics. For a description and comparison of the methods of expansion used in these two fields, see the Appendix of Ref. 7. We have checked our series against those in the literature wherever comparison is possible, and review the most stringent checks below. In all cases, our series agreed with those found elsewhere.

The most general work preceding ours is that of Kincaid *et al.*<sup>(3)</sup> Their series for the susceptibility  $\chi$  and second moment  $\mu_2$  are identical to ours through tenth order (the highest computed by them) for all four of the lattice geometries we considered.

To check our results at an order higher than ten, we must rely on special choices of the potential. The spin- $S$  Ising model has been well studied. The moments  $I_{2n}$  in this model can be easily derived, since there is no self-interaction term [ $V(\phi)=0$ ], and the off-diagonal nearest neighbor spin-spin interaction plays the role of the kinetic energy operator in lattice field theory. The internal line counter and expansion parameter  $v$  is replaced with the coupling parameter  $J/kTS^2$ , where  $J$  is now the spin exchange constant and the external field is the magnetic field  $H$ . With these identifications and Eqs. (5.1) and (2.5), the moments  $I_{2n}$  in the spin- $S$  Ising model can be found from

$$\sum_{n=0}^{\infty} \frac{I_{2n} h^{2n}}{(2n)!} = \frac{\sinh[Y(h/2)]}{Y \sinh(h/2)} \quad (6.1)$$

(where  $Y=2S+1$  and  $h=H/kT$ ), by expanding the left-hand side and matching like powers of  $h$ .

For example, for the spin-1/2 Ising model,  $Y=2$  and one finds from Eq. (6.1) that

$$I_{2n} = 1/(2)^{2n}, \quad v = 4J/kT \quad (6.2)$$

With these substitutions, the linear chain (LC) spin-1/2 results are trivial and easily checked against closed-form expressions. Our plane square (PSQ) series for  $\chi$  and  $\mu_2$  agree with those of Wu *et al.*<sup>(8)</sup> through 15th order. Our series on a simple cubic (SC) lattice match those of Sykes *et al.*<sup>(9)</sup> for  $\chi$  through 15th order, those of Moore *et al.*<sup>(10)</sup> for  $\mu_2$  through order 12, and, of course, those of Roskies<sup>(4)</sup> for  $\mu_2$  through 15 orders. The most extensive spin-1/2 series for  $\chi$  and  $\mu_2$  on a body-centered cubic (BCC) lattice are the 21st-order results of Nickel,<sup>(11)</sup> which agree through order 15 with our own.

In conclusion, we have extended through 15th order the strong coupling (high-temperature) series for the susceptibility  $\chi$  and second moment  $\mu_2$  for models that are even in a scalar potential on four different lattice geometries: linear chain (LC), plane square (PSQ), simple cubic (SC), and body-centered cubic (BCC). Our algorithm for generating these series has the advantage of separating lattice-dependent from model-dependent considerations, and is computationally efficient. The graph enumeration scheme we use is new, and is described in the Appendix. Less general series in the literature are consistent with our own.

## APPENDIX. GENERATION OF GRAPHS

The basic problem in a graphical enumeration scheme is to ensure that each different graph is counted once and only once. If one has an algorithm for generating graphs, one must generally develop a canonical form for graphs and accept a new graph only if it is in canonical form. One has also to assure that every canonical form is generated once and only once. This is a difficult problem (see, e.g., Ref. 12). It sometimes requires searching through a large list of already generated graphs to see that the newly generated one is different from all of those in the list.

But in our case, as explained in the text, we multiply the contribution of each graph by its topological symmetry number. So we can afford to generate a graph more than once, provided that we associate with it a pseudosymmetry number, which has the property that the sum of the pseudo-symmetry numbers for the graphs equivalent to a given graph is the correct symmetry number for that graph. This is usually wasteful, but we were able to use it to find a graph-generating algorithm that eliminated the need for keeping track of the graphs already encountered.

It is easy to find a canonical representation for a graph. Represent it by a symmetric connection matrix whose rows and columns represent the vertices and whose  $(i, j)$ th entry indicates the number of lines joining vertex  $i$  to vertex  $j$ . We can introduce an ordering of matrices (say lexical ordering reading row by row). A graph will be in canonical form if its corresponding matrix is the highest of all equivalent graphs obtained by an arbitrary permutation of its vertices. The problem then is to check whether a given matrix is lexically higher than all equivalent matrices obtained by an arbitrary permutation of its vertices. This is potentially very time-consuming, because the number of permutations grows rapidly with the size of the graph. There are trivial simplifications, such as ordering the vertices by their degree and considering only subgroups of permutations that permute only vertices of the same degree. This is still too time-consuming to be practical.

We present a graph-generating scheme that will sometimes generate copies of equivalent graphs, but will associate to them a pseudo-symmetry number that has the desired property mentioned above. As the degree of the graphs grows, the number of copies of equivalent graphs grows, too, but to the order we work (15 lines), this overcounting is not serious. It would be more severe if we went to much higher order (as Nickel<sup>(11)</sup> has done).

Suppose that a particular  $N \times N$  connection matrix  $M$  is given. We label its vertices 1 through  $N$ , corresponding to its rows and columns. A matrix obtained by permuting the rows and column of  $M$  clearly represents the same graph, but may be represented by a different matrix. For example, the graph shown in Fig. 5 with two different vertex labelings can be represented by two different corresponding matrices.

What follows is the algorithm for accepting a particular connection matrix and computing its pseudo-symmetry number  $ps$ . First set  $ps = 1$ . Then examine the vertices in order, making sure that they have been chosen to satisfy certain criteria. At any stage we assume that  $NV$  vertices have already been chosen, and that they are numbered 1 to  $NV$ . The next vertex must satisfy the following conditions:

1. It must be the highest degree vertex remaining. If there are several, it must be the one connected to the lowest-numbered chosen vertex. If there are several, it must be the one also connected to the next-lowest-numbered chosen vertex, and so on. If there are still several, pick the one connected to the highest number of unchosen neighbors of the lowest-numbered chosen vertex. If there are several, it must be the one connected to

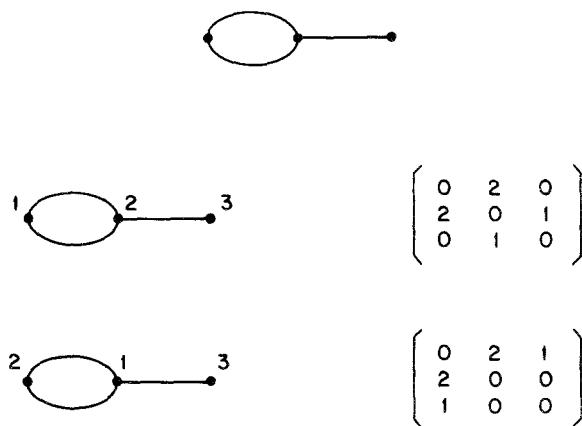


Fig. 5. Connection matrices corresponding to two vertex labelings of a simple graph.

the highest number of unchosen neighbors of the next-lowest-numbered chosen vertex, etc. (The distinguishing feature of this algorithm is that it stops after examining the chosen vertices or their nearest neighbors only.)

2. If there are still several candidate vertices, divide  $ps$  by the number of such candidate vertices.

3. If there are still unchosen vertices, go to step 1. If not, the graph matrix is acceptable and its pseudo-symmetry number is the value of  $ps$ .

As an example, consider the graph shown in Fig. 6. Its first vertex is unique (the one of degree 4). The next vertex can be chosen in four

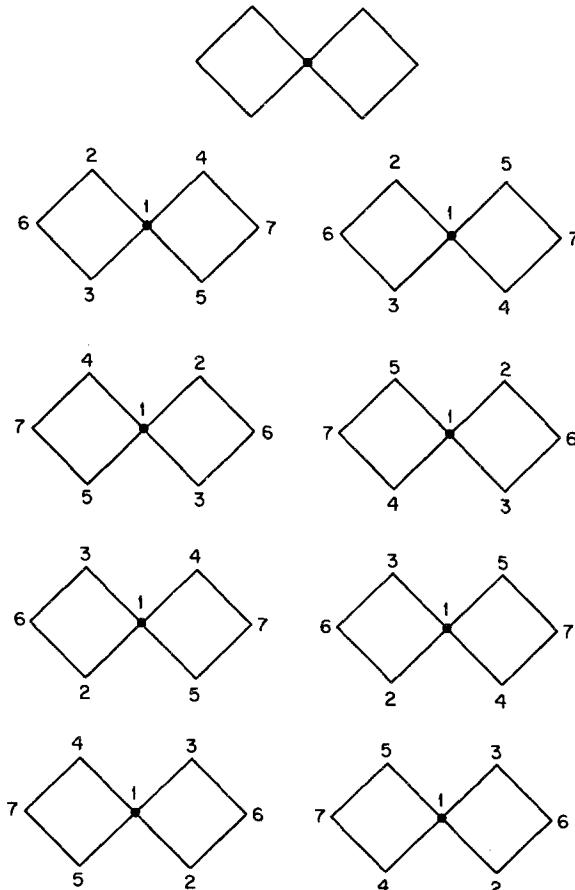


Fig. 6. Vertex labelings consistent with the algorithm given in the text. Since there are eight such labelings, the connection matrix associated with these labelings would have associated with it a pseudo-symmetry number of 1/8.

equivalent ways. The position of the third vertex is then uniquely determined by the rules of step 1. Having chosen vertices 2 and 3 as shown in the first column of the figure, we can choose the fourth vertex in two ways. The choices for the remaining vertices are completely determined by the algorithm. The pseudo-symmetry number for this graph is  $1/8$ , which coincides with the true symmetry number.

As a second example, consider the graph shown in Fig. 7. It will appear as three different acceptable connection matrices, corresponding to the vertex labelings shown in Fig. 7 are  $1/8$ ,  $1/24$ ,  $1/12$ . The true symmetry number is  $1/4 = 1/8 + 1/24 + 1/12$ .

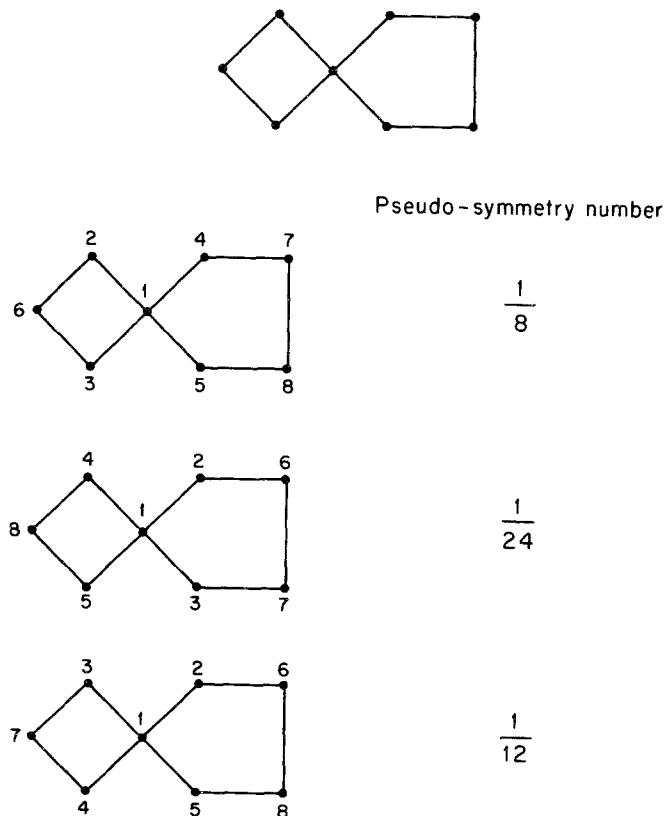


Fig. 7. An example of labelings for each of three acceptable connection matrices associated with the graph shown. According to the algorithm given in the text, each matrix would be accepted with the pseudo-symmetry number shown, which sum to the true symmetry number of the graph.

**Table I.** Strong Coupling Expansion of the Susceptibility  $\chi$  through Fifteenth Order<sup>a</sup>

Partition	Linear Chain (LC)	Plane Square (PSQ)	Simple Cubic (SC)	Body-Centered Cubic (BCC)
- 0 -	( 1, 0, 0, 0, 0, 0, 0, 0 )	1	1	1
- 1 -	( 2, 0, 0, 0, 0, 0, 0, 0 )	2	4	6
- 2 -	{ 3, 0, 0, 0, 0, 0, 0, 0 } ( 1, 1, 0, 0, 0, 0, 0, 0 )	2	20	104
- 3 -	{ 4, 0, 0, 0, 0, 0, 0, 0 } ( 2, 1, 0, 0, 0, 0, 0, 0 ) ( 0, 2, 0, 0, 0, 0, 0, 0 )	2	4	54
- 4 -	{ 5, 0, 0, 0, 0, 0, 0, 0 } ( 3, 1, 0, 0, 0, 0, 0, 0 ) ( 2, 0, 1, 0, 0, 0, 0, 0 ) ( 1, 2, 0, 0, 0, 0, 0, 0 ) ( 0, 1, 1, 0, 0, 0, 0, 0 )	-6 -6 -6 -6 -6	132 72 4 164 2	1992 180 16 414 6
- 5 -	{ 6, 0, 0, 0, 0, 0, 0, 0 } ( 4, 1, 0, 0, 0, 0, 0, 0 ) ( 3, 0, 1, 0, 0, 0, 0, 0 ) ( 2, 2, 0, 0, 0, 0, 0, 0 ) ( 1, 1, 1, 0, 0, 0, 0, 0 ) ( 0, 0, 2, 0, 0, 0, 0, 0 )	240 -40 0 210 60 2	10560 11760 5460 3600 22960 360	248400 127400 3600 22960 900 4
- 6 -	{ 7, 0, 0, 0, 0, 0, 0, 0 } ( 5, 1, 0, 0, 0, 0, 0, 0 ) ( 4, 0, 1, 0, 0, 0, 0, 0 )	720 -360 -360	132480 154440 11160	6162480 3831200 128880

{ 3, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	-780	116760	10135800	4321200			
{ 3, 0, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	360	14400	5040			
{ 2, 1, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	330	95400	64440	178560			
{ 1, 3, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	30	180	37710	109320			
{ 1, 1, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	52	320	804	840			
{ 0, 2, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	70	420	1050	1504			
{ 0, 0, 1, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	2	4	6	1960			
{ 0, 0, 0, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }				a			
<hr/>							
{ 6, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	-16380	1917720	178230780	2191185360			
{ 6, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	40320	2439360	135898560	12433920			
{ 5, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	100800	44963100	25100400			
{ 4, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	-28140	2221800	44983100	44915040			
{ 4, 0, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	5040	75600	75600			
{ 3, 1, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	-4200	438480	3767400	16070880			
{ 2, 3, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	5040	537600	3989720	16813440			
{ 2, 1, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	10080	50400	141120			
{ 2, 0, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	924	20552	85092	2080832			
{ 1, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	2240	47040	188160	519680			
{ 1, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	2112	672	1680	31386			
{ 0, 4, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	70	12460	43890	151480			
{ 0, 2, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	140	840	2100	3920			
{ 0, 0, 0, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	2	4	6	b			
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{ 9, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	-65520	29161440	5855118320	10376644160			
{ 7, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	63000	47612880	5174946520	7067204480			
{ 6, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	37800	1154160	17462440	190831520			
{ 5, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	63000	41277600	2039648280	2035614400			
{ 5, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	17640	3107160	27634720			
{ 4, 1, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	-59640	9176160	189108360	129121600			
{ 4, 0, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	19078080	303516360	1939677600			
{ 3, 3, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	-63360	3899840	3931200	17424960			
{ 3, 1, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	-3360	634992	5668992	23816576			
{ 3, 0, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	-1792	11480	3461320	26412120			
{ 2, 2, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	5040	50400	2521200			
{ 2, 1, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	1400	46592	206136	70560			
{ 2, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	11080	1246280	8921100	566720			
{ 1, 4, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	1756	754660	892100	41637920			
{ 1, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	5264	12080	342290	342290			
{ 1, 1, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	56	336	46818	1392388			
{ 1, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	86	524	840	1588			
{ 1, 0, 0, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	1540	83720	1314	2456			
{ 0, 3, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	70	420	293580	980560			
{ 0, 2, 0, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	420	2520	1050	1980			
{ 0, 1, 1, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	2	4	6	11760			
{ 0, 0, 0, 1, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }				b			

- 7 -

- 8 -

<sup>a</sup> The results are completely general for any even potential in zero external field for the lattice geometries linear chain (LC), plane square (PSQ), simple cubic (SC), and body-centered cubic (BCC). The partitions are grouped by perturbation order and the corresponding coefficients in each *n*th order have been multiplied by *n!* for convenience.

(Table continued)

Table I. (*Continued*)

1  
8

10

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(Table continued)

Table I. (*Continued*)

6

Table I. (*Continued*)

(Table continued)

**Table I.** (*Continued*)

4, 0, 2, 0, 1, 0, 0, 0, 0)	-22500508784	636420698912	2264550557929940
4, 0, 1, 2, 0, 0, 0, 0, 0)	-718377660	770021932680	2557245528540
4, 0, 1, 0, 0, 0, 0, 0, 0)	0	17797120	22726200
4, 0, 0, 1, 0, 0, 0, 0, 0)	-347444	7270569316400	58373200
4, 0, 0, 1, 0, 0, 0, 0, 0)	-597417890	7270569316400	734381981406600
3, 6, 0, 0, 0, 0, 0, 0, 0)	-57322517200	8514207200	1842089904493200
3, 4, 0, 1, 0, 0, 0, 0, 0)	-63463995060	3466966742036800	10490056458106000
3, 3, 2, 0, 0, 0, 0, 0, 0)	0	148481199500	1049381781890400
3, 2, 1, 0, 0, 0, 0, 0, 0)	-104951645520	459551645520	1842089915839580
3, 2, 0, 2, 0, 0, 0, 0, 0)	-1911259350	4595518377280	1947599915839580
3, 2, 0, 0, 0, 0, 0, 0, 0)	-5322517200	1266643057600	162124505691200
3, 1, 2, 1, 0, 0, 0, 0, 0)	0	1266643057600	1842089915839580
3, 1, 1, 0, 0, 0, 0, 0, 0)	-24504480	25415319760	10490056458106000
3, 1, 0, 1, 0, 0, 0, 0, 0)	-29129100	23068913280	10490056458106000
3, 0, 4, 0, 0, 0, 0, 0, 0)	-44972528	203989512064	10490056458106000
3, 0, 2, 0, 0, 0, 0, 0, 0)	-18866648	203989594916	10490056458106000
3, 0, 1, 1, 0, 0, 0, 0, 0)	-62486474	7301251312	10490056458106000
3, 0, 0, 3, 0, 0, 0, 0, 0)	0	804840000	11853475560
3, 0, 0, 1, 0, 0, 0, 0, 0)	0	1081000	16216200
3, 0, 0, 0, 1, 0, 0, 0, 0)	0	-48048	1899880
3, 0, 0, 0, 0, 2, 0, 0, 0)	-10039679560	21961682124400	18254551961554000
3, 0, 0, 0, 0, 1, 0, 0, 0)	-1898980	184099635700	6350562450200
2, 3, 1, 0, 0, 0, 0, 0, 0)	-2583060480	2552753901140	6596138110242720
2, 3, 0, 0, 0, 0, 0, 0, 0)	-234594380	1931929870400	47325205333960
2, 2, 3, 0, 0, 0, 0, 0, 0)	-12612600	89929162860	161439189980
2, 2, 2, 0, 0, 0, 0, 0, 0)	-12432420	27847089160	390813934500
2, 2, 1, 2, 0, 0, 0, 0, 0)	-6040448	32776693884	4484911406976
2, 1, 1, 2, 0, 0, 0, 0, 0)	2589188650	516008493000	684014713000
2, 1, 1, 0, 0, 0, 0, 0, 0)	0	15135120	227026800
2, 1, 0, 1, 0, 0, 0, 0, 0)	-360360	29595200	28313485120
2, 1, 0, 0, 1, 0, 0, 0, 0)	0	8605910552	6671444
2, 0, 3, 1, 0, 0, 0, 0, 0)	0	30605910552	3913663041288
2, 0, 2, 0, 0, 0, 0, 0, 0)	0	266378112	25574353520
2, 0, 1, 0, 0, 0, 0, 0, 0)	0	1829676240	1399658976
2, 0, 0, 2, 0, 0, 0, 0, 0)	0	705727268	51050400
2, 0, 0, 1, 0, 0, 0, 0, 0)	0	536720184	3816930400
2, 0, 0, 0, 1, 0, 0, 0, 0)	0	12612000	18918900
2, 0, 0, 0, 0, 1, 0, 0, 0)	0	134787421080	1613197114600
1, 709007300	18705557871000	73955211703900	13732712368174800
1, 709007300	375789036200	1124705200	1315200000000000
1, 709007300	181059348580	5759408613198740	10390160680
1, 709007300	1715353800	34481574000	955772989712
1, 709007300	41261093800	53287760766240	3412891303044
1, 709007300	3204525220	3513694298160	13499130544
1, 709007300	12612000	88288200	1151221616
1, 709007300	813333120	165159327340	61994145224
1, 709007300	31735050	38324304000	12788519824
1, 709007300	4144140	233655820	10390160680
1, 709007300	12292200	1999392552	45676040
1, 709007300	5586340956	3280810392	45676040
1, 709007300	3897884	2895312988	15135120
1, 709007300	91323132	1151616125	16746300
1, 709007300	185858450	205619520	1101000000000000
1, 709007300	10810580	5054000	15135120
1, 709007300	0	5989984	5989984
1, 709007300	0	4101582	4101582
1, 709007300	0	152152	152152
1, 709007300	0	198016	198016
1, 709007300	0	1676696	1676696

(Table continued)

Table I. (Continued)

{ 7, 3, 1, 0, 0, 0, 0, 0 }	-38202987824000	11824041581520000	1593121174042477456000	957717801043531296000
{ 7, 3, 0, 1, 0, 0, 0, 0 }	-117145288000	1552021853376000	210947608991472000	11394608087896032000
{ 7, 1, 0, 0, 0, 0, 1, 0 }	0	19875015006400	53675217840470000	960613671390497600
{ 7, 1, 0, 0, 0, 0, 0, 0 }	0	32891592000	12259447400	52390287904000
{ 7, 0, 1, 0, 0, 1, 0, 0 }	0	1100755656000	372546255476400	24685403945435937600
{ 7, 0, 1, 0, 0, 0, 0, 0 }	0	147130547788226000	3205402026364160000	1892185586233600
{ 6, 5, 0, 0, 0, 0, 0, 0, 0 }	0	26339927788226000	1381052026364160000	100438779561104000
{ 6, 3, 0, 1, 0, 0, 0, 0, 0 }	0	698809713600	46156168821085230000	778419498182015286000
{ 6, 2, 0, 2, 0, 0, 0, 0, 0 }	0	19978791896860800	14269120459788176000	2746655518717192000
{ 6, 2, 0, 1, 0, 0, 0, 0, 0 }	0	1225944720000	2102495194800000	6146704133716264000
{ 6, 1, 0, 1, 0, 0, 0, 0, 0 }	0	175204149120000	55796005220152000	11328822373688000
{ 6, 1, 0, 0, 2, 0, 0, 0, 0 }	0	136442133828000	56619245586102000	187448818760934400
{ 6, 1, 0, 0, 0, 0, 0, 0, 0 }	0	0	0	1671057521891514000
{ 6, 0, 2, 1, 0, 0, 0, 0, 0 }	0	16710027302934400	70853669866432000	152562098000
{ 6, 0, 1, 1, 0, 0, 0, 0, 0 }	0	0	25714839110000	21492447096543400
{ 6, 0, 0, 1, 1, 0, 0, 0, 0 }	0	0	572243861670000	1156168772277000
{ 6, 0, 0, 0, 1, 1, 0, 0, 0 }	0	0	7935011534400	126148993094400
{ 6, 0, 0, 0, 0, 1, 1, 0, 0 }	0	0	138577295834400	138577295834400
{ 5, 7, 1, 0, 0, 0, 0, 0, 0 }	0	71143605261725000	3111895776938824000	128137408798816000
{ 5, 7, 1, 0, 0, 0, 0, 0, 0 }	0	40216762377792000	8908730378640000	307321257518912000
{ 5, 2, 1, 1, 0, 0, 0, 0, 0 }	0	99046720517500000	99046720517500000	29715538062818960000
{ 5, 2, 1, 0, 0, 0, 0, 0, 0 }	0	0	15038253232000	185125042000000
{ 5, 1, 3, 0, 0, 0, 0, 0, 0 }	0	3019065953904000	555701778099380000	163183102531666500
{ 5, 1, 1, 0, 1, 0, 0, 0, 0 }	0	0	674122840000	2128007740982400
{ 5, 1, 0, 1, 1, 0, 0, 0, 0 }	0	30728995300000	3543380782800000	6586743000005600
{ 5, 0, 2, 0, 1, 0, 0, 0, 0 }	0	10829628800	3346718377883600	4385718388066560
{ 5, 0, 1, 1, 0, 0, 0, 0, 0 }	0	1324828048000	28369713883600	5227331031724800
{ 5, 0, 0, 1, 1, 0, 0, 0, 0 }	0	0	2824456114520000	305124019200
{ 5, 0, 0, 0, 1, 0, 0, 0, 0 }	0	0	1089728400	1089728400
{ 5, 0, 0, 0, 0, 1, 0, 0, 0 }	0	0	54116702400	54116702400
{ 5, 0, 0, 0, 0, 0, 1, 0, 0 }	0	0	3156217383320	3543434995692994400
{ 5, 0, 0, 0, 0, 0, 0, 1, 0 }	0	0	11775301531456666000	457283434995692994400
{ 4, 6, 1, 0, 0, 0, 0, 0, 0 }	0	4393083312363600	798061054399710000	2249898911114944000
{ 4, 6, 1, 0, 0, 0, 0, 0, 0 }	0	42033770590000	2630793050518200	6000597693793187000
{ 4, 6, 1, 0, 0, 0, 0, 0, 0 }	0	213758167489600	2125146794855000	25562778895393000
{ 4, 6, 1, 0, 0, 0, 0, 0, 0 }	0	0	121571572956000	47432114618901600
{ 4, 5, 0, 1, 0, 0, 0, 0, 0 }	0	220433635925600	190303348556400	35531026172716800
{ 4, 5, 0, 1, 0, 0, 0, 0, 0 }	0	2742627709800	6610304000	1025392145979801600
{ 4, 5, 0, 0, 1, 0, 0, 0, 0 }	0	0	51601343187794400	182654571241600
{ 4, 4, 1, 2, 1, 0, 0, 0, 0 }	0	756284995065600	9020228817600	1945944000
{ 4, 4, 1, 1, 2, 0, 0, 0, 0 }	0	-1362160000	62891838810800	1344527096000
{ 4, 4, 1, 0, 1, 0, 0, 0, 0 }	0	134363108800	5125424745750	722839763500
{ 4, 4, 1, 0, 0, 1, 0, 0, 0 }	0	147471209960	77121770049040	1599073342018760
{ 4, 4, 0, 1, 0, 0, 1, 0, 0 }	0	13526375590280	491836287040	5629768972440
{ 4, 4, 0, 1, 0, 0, 0, 1, 0 }	0	0	49090903207200	1629873018600
{ 4, 4, 0, 0, 1, 0, 1, 0, 0 }	0	47844276800	47844276800	208487172654400
{ 4, 4, 0, 0, 0, 1, 0, 0, 1 }	0	0	1163242080	18162144000
{ 4, 4, 0, 0, 0, 0, 1, 0, 1 }	0	0	29072403360	208487443840
{ 4, 0, 0, 0, 0, 0, 2, 0, 0 }	0	-6455180	3348519340	21557682960

(Table continued)

Table I. (*Continued*)

Table II. Strong Coupling Expansion of the Second Moment  $\mu_2$  through Fifteenth Order

Partition	Linear Chain (LC)	Plane Square (PSQ)	Simple Cubic (SC)	Body-Centered Cubic (BCC)
- 0 -		0	0	0
( 1, 0, 0, 0, 0, 0, 0, 0 )	0			
- 1 -				
{ 2, 0, 0, 0, 0, 0, 0, 0 }	2	4		8
- 2 -				
( 3, 0, 0, 0, 0, 0, 0, 0 )	16	64	144	256
- 3 -				
{ 4, 0, 0, 0, 0, 0, 0, 0 }	90	900	3294	8136
{ 2, 1, 0, 0, 0, 0, 0, 0 }	12	72	180	336
{ 0, 2, 0, 0, 0, 0, 0, 0 }	2	4	6	8
- 4 -				
{ 5, 0, 0, 0, 0, 0, 0, 0 }	384	13056	82944	291840
{ 3, 1, 0, 0, 0, 0, 0, 0 }	192	3012	12046	30720
{ 1, 2, 0, 0, 0, 0, 0, 0 }	64	256	576	1024
- 5 -				
{ 6, 0, 0, 0, 0, 0, 0, 0 }	1200	202560	2347920	11847360
{ 4, 1, 0, 0, 0, 0, 0, 0 }	1440	88560	594000	2155680
{ 3, 0, 1, 0, 0, 0, 0, 0 }	0	720	3600	10080
{ 2, 2, 0, 0, 0, 0, 0, 0 }	1170	13140	4890	12420
{ 1, 1, 1, 0, 0, 0, 0, 0 }	60	360	900	1680
{ 0, 0, 12, 0, 0, 0, 0, 0 }	2	4	6	8
- 6 -				
{ 7, 0, 0, 0, 0, 0, 0, 0 }	4320	3412800	74766240	542056320
{ 5, 1, 0, 0, 0, 0, 0, 0 }	1440	223520	2784480	13708480
{ 4, 0, 1, 0, 0, 0, 0, 0 }	0	40320	336960	1313280
{ 3, 2, 0, 0, 0, 0, 0, 0 }	14880	575040	1307040	1330720
{ 2, 1, 1, 0, 0, 0, 0, 0 }	1440	26880	108000	276480
{ 1, 3, 0, 0, 0, 0, 0, 0 }	960	14400	51840	136320

(	1.0,-2,0,0,0,0,0,0)	96	384	984
(	0,2,1,0,0,0,0,0)	160	640	1440

(	1.0,-2,0,0,0,0,0,0)	)
(	0,2,1,0,0,0,0,0)	)

- 7 -

(	B,0,0,0,0,0,0,0,0)	44100	62881560	2655431100	27689553360
(	6,1,0,0,0,0,0,0,0)	-40640	5485360	120059200	981566240
(	5,0,1,0,0,0,0,0,0)	0	1391040	21833280	12605640
(	4,2,0,0,0,0,0,0,0)	116340	21602280	24696540	1311379440
(	4,0,1,0,0,0,0,0,0)	0	5040	75600	352800
(	3,1,1,0,0,0,0,0,0)	15960	1459040	10117800	38004960
(	3,1,1,0,0,0,0,0,0)	31920	1559040	9802800	37242240
(	2,3,0,0,0,0,0,0,0)	0	10080	50400	141120
(	2,1,0,0,0,0,0,0,0)	0	36680	139524	359856
(	2,0,2,0,0,0,0,0,0)	2940	36680	309120	806400
(	1,2,1,0,0,0,0,0,0)	6720	82890	1680	3136
(	1,0,1,0,0,0,0,0,0)	112	672	104370	294940
(	1,0,0,0,0,0,0,0,0)	2310	30380	2100	3920
(	0,2,0,1,0,0,0,0,0)	140	840	6	8
(	0,0,0,2,0,0,0,0,0)	2			

- 8 -

(	9,0,0,0,0,0,0,0,0)	322560	1264112640	104309130240	1565864294400
(	7,1,0,0,0,0,0,0,0)	-483840	1363944960	62575027200	701510968800
(	6,0,1,0,0,0,0,0,0)	0	40158720	1228719120	1105057600
(	5,2,0,0,0,0,0,0,0)	120960	73340800	1593160000	1249442400
(	5,0,0,0,0,0,0,0,0)	0	312560	8709120	58000800
(	4,1,0,0,0,0,0,0,0)	0	6492080	818657280	455068720
(	4,0,1,0,0,0,0,0,0)	0	10898400	1208511360	667724000
(	3,3,0,0,0,0,0,0,0)	470400	8086400	1205680	27740160
(	3,1,0,1,0,0,0,0,0)	0	43904	245376	54622240
(	3,0,1,0,0,0,0,0,0)	0	43904	8798720	216933680
(	2,2,1,0,0,0,0,0,0)	170240	3584	68096	9503520
(	2,0,1,0,0,0,0,0,0)	0	4067840	2378800	1361920
(	1,4,0,0,0,0,0,0,0)	116480	116480	524160	157376
(	1,2,0,1,0,0,0,0,0)	0	4480	16656	604800
(	1,1,2,0,0,0,0,0,0)	11648	11648	1152	2048
(	1,0,0,2,0,0,0,0,0)	0	4480	107512	362880
(	0,3,0,0,0,0,0,0,0)	0	4480	107512	8064
(	0,1,1,0,0,0,0,0,0)	896	3584	0	14336

- 9 -

(	10,0,0,0,0,0,0,0,0)	27569263680	4495820565600	9720332016000
(	8,-1,0,0,0,0,0,0,0)	5443200	35322557760	5262866348800
(	7,0,1,0,0,0,0,0,0)	0	1079385560	7388912880
(	6,2,0,0,0,0,0,0,0)	-6342840	238812785880	1027211956240
(	6,0,1,0,0,0,0,0,0)	0	10535520	683788240
(	5,1,0,1,0,0,0,0,0)	-861840	2521683360	6107662800
(	5,0,0,1,0,0,0,0,0)	0	0	12061386800
(	4,3,0,0,0,0,0,0,0)	710640	5821653600	120716726480
(	4,1,0,1,0,0,0,0,0)	0	42910560	726933360
(	4,0,2,0,0,0,0,0,0)	410760	114402960	135302840
(	3,2,0,0,0,0,0,0,0)	985760	649831440	7651040160
(	3,1,0,0,1,0,0,0,0)	0	0	151200

(Table continued)

Table II. (*Continued*)

11

384199200	(12.0,0.0,0.0,0.0,0.0,0.)	107366978233600	4818978225606400
-1766318400	(10.1,0.0,0.0,0.0,0.0,0.)	27764617657600	3526140745184000
0	(9.0,1.0,0.0,0.0,0.0,0.)	83661717124600	75854772448000
2799165600	(8.2,0.0,0.0,0.0,0.0,0.)	255543938400	45776688912000
0	(8.0,0.1,0.0,0.0,0.0,0.)	5019337600	298223212000
133056000	(7.1,1.0,0.0,0.0,0.0,0.)	305995536000	74755921502000
0	(7.0,0.1,0.0,0.0,0.0,0.)	0	58277177400
-1570060800	(6.3,0.0,0.0,0.0,0.0,0.)	114581340459200	3539315200
0	(6.1,0.1,0.0,0.0,0.0,0.)	68866459200	162879328000
5544000	(6.0,2.0,0.0,0.0,0.0,0.)	193258456320	5230533977600
0	(5.9,0.0,0.0,0.0,0.0,0.)	0	869177981600
-358696800	(5.2,1.0,0.0,0.0,0.0,0.)	1974462336000	10980617199520
0	(5.1,0.0,1.0,0.0,0.0,0.)	488880800	41912400
-2772000	(5.0,1.0,0.1,0.0,0.0,0.)	16705643200	112145135645000
0	(5.0,0.1,1.0,0.0,0.0,0.)	0	37901001600
-2403818040200	(4.4,0.0,0.0,0.0,0.0,0.)	24047100	3660594981120
-3456000	(4.2,0.1,0.0,0.0,0.0,0.)	55519002000	108374799400
-1774080	(4.1,2.0,0.0,0.0,0.0,0.)	129286080000	147390813000
0	(4.1,0.0,0.0,0.0,0.0,0.)	0	286901351360
1047420	(4.1,0.0,1.0,0.0,0.0,0.)	2244199360	74844000
139154400	(4.0,0.2,0.0,0.0,0.0,0.)	498471880	3951759640
0	(4.0,0.1,0.0,0.0,0.0,0.)	0	6140805080
20180160	(3.9,0.0,0.0,0.0,0.0,0.)	388678460000	703002448000
0	(3.8,-1.0,0.0,0.0,0.0,0.)	8709345760	840987984000
0	(3.0,-3.0,0.0,0.0,0.0,0.)	1199056320	61612290320
1306800	(3.0,0.1,0.0,0.0,0.0,0.)	17178480	95936174720
0	(2.5,0.0,1.0,0.0,0.0,0.)	120195864000	1520109120
172557000	(2.3,0.1,0.0,0.0,0.0,0.)	2216979720	2160517174200
93970800	(2.2,0.0,0.0,0.0,0.0,0.)	2216979720	7520835400
0	(2.2,0.0,0.0,0.0,0.0,0.)	0	236877102000
6098940	(2.1,1.0,0.0,0.0,0.0,0.)	2494800	37122000
151920	(2.1,0.2,0.0,0.0,0.0,0.)	143146380	103567440
1413720	(2.0,2.1,0.0,0.0,0.0,0.)	105407280	69922280
0	(2.0,0.1,0.0,0.0,0.0,0.)	91153200	596340360
11484	(1.9,1.0,0.0,0.0,0.0,0.)	71880	356400
82982800	(1.8,1.0,0.0,0.0,0.0,0.)	12135261600	9909035600
0	(1.7,0.0,0.0,0.0,0.0,0.)	0	11990935600
5313000	(1.6,1.0,0.0,0.0,0.0,0.)	52668000	384199200
6244720	(1.5,1.0,0.0,0.0,0.0,0.)	55717200	3486133520
5930800	(1.4,1.0,0.0,0.0,0.0,0.)	40482980	2379540240
162360	(1.3,1.0,0.0,0.0,0.0,0.)	665280	322640
155232	(1.2,1.0,0.0,0.0,0.0,0.)	228880	8155860
264	(1.1,0.0,0.0,0.0,0.0,0.)	0	8034988
1313466000	(1.0,0.0,0.0,0.0,0.0,0.)	1313466000	11684395800
0	(0.9,0.0,0.0,0.0,0.0,0.)	0	84866628000
1386000	(0.8,0.0,0.0,0.0,0.0,0.)	149985200	43220000
2079000	(0.7,0.0,0.0,0.0,0.0,0.)	117556200	67443400
0	(0.6,0.0,0.0,0.0,0.0,0.)	0	2776200
280170	(0.5,0.0,0.0,0.0,0.0,0.)	476580	1107577280
388080	(0.4,0.0,0.0,0.0,0.0,0.)	7318080	24181680
1980	(0.3,0.0,0.0,0.0,0.0,0.)	11880	52440
47124	(0.2,0.0,0.0,0.0,0.0,0.)	212332	24049872
1848	(0.1,0.0,0.0,0.0,0.0,0.)	0	6627852
	(0.0,0.0,0.0,0.0,0.0,0.)	0	51774

(Table continued)

Table II. (Continued)

- 12 -

{ 13, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	- 5748019200	58780081867334400	3807308780829491200
{ 11, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	- 14840409600	633457567556354000	3144244987288662800
{ 10, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	17289689687424000	7160334301952000
{ 9, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	3677345600	32 18775848903617200	1177439768457535000
{ 9, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	99633333800	20147831994600	7689231389388800
{ 8, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	10537113295200	2323893558759800	63970391607912000
{ 8, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	17403724800	105324518400	4075337213400
{ 7, 3, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	475170743536900	81619577434400	21891648329566400
{ 7, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	2326275393000	41619577434400	888990106896400
{ 7, 1, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	41619577434400	12866924589600	12866924589600
{ 7, 0, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	149022720	69291882874680	91988307200
{ 7, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	215507200	160313435
{ 6, 5, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	12773331600	9000125400400	12612600
{ 6, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	1460548800	617933456000
{ 6, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	- 26611200	3528327786600	4333800
{ 6, 4, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	137868688742400	1021650681300000	18729959878656000
{ 6, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	- 32265000	15859835348000	188229598724400
{ 5, 5, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	- 968647680	322480305532000	43161413170760
{ 5, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	12211054800	17440576000
{ 5, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	447188686400	4225597080
{ 4, 5, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	661008430800	56193548160
{ 4, 3, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	- 522910800	112586176844800	148748612993200
{ 4, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	170017826400	1665231577600
{ 4, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	65520000	144261232960	12906666074560
{ 4, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	- 7883600	265793197240	24906658197120
{ 4, 0, 0, 3, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	1868160240	127119720
{ 4, 0, 0, 0, 3, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	17886710800	10254598080
{ 3, 5, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	- 6747840	137009856000	59880379672000
{ 3, 3, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	2840838392000	62948395800	14270837852800
{ 3, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	1130285319200	440380787560
{ 3, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	6192852319200	7666026760	4142595084800
{ 3, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	109137854400	179713419960	537462400
{ 3, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	956400	119280909240	1072906803200
{ 3, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	143528595640	85024534400
{ 3, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	216160000	141884266600	8572980800
{ 3, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	1191915480	35664640
{ 3, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	97366600	59982400
{ 3, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	17988640	123111816
{ 2, 4, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	2621203200	2081971584000	38248317678800
{ 2, 3, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	7460508400	10957161600
{ 2, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	194696860	9161171720	103934003200
{ 2, 1, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	183391360	549201945600	584065278720
{ 2, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	1463160	940282400	886154960
{ 2, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	416526460	3591837420
{ 2, 0, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	33015288	1083615040
{ 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	16422120	43003636920
{ 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	2081971584000	3617245382400
{ 1, 6, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	561488000	292967136000	7056934200
{ 1, 6, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	36092800	25591654400	67371524800
{ 1, 3, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	392071680	81531994240	402481626240
{ 1, 3, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	53222400	3591837420
{ 1, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	2211600	751322880	2797022400
{ 1, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	1018195200	1684812800
{ 1, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	24992	108003560
{ 1, 2, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	9881416	43003636920
{ 1, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	561488000	4907311516800	48553491571200
{ 1, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	36092800	327038344200	230738188800
{ 1, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	392071680	8615897160	601147987200
{ 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	56883200	2341785600
{ 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	5030843360	2126348600
{ 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	8454548624	26549972400
{ 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	1624472680	124472680
{ 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	12266160	5800475220
{ 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	6747840	31933440
{ 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	12317160	11277160
{ 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	13539188988	13539188988
{ 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	11970400	31283808
{ 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	5601024	52867584
{ 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	392832	9982127200
{ 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	192	518911400
{ 0, 5, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	3769200	15552424400	2483712000
{ 0, 4, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	17740400	41003085120
{ 0, 3, 1, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0, 0 }	0	24097920	5275670400

{ 0, 2, 1, 0, 0, 1, 0, 0 }	212889960	56770560
{ 0, 1, 0, 1, 0, 0, 0, 0 }	47308880	78566400
{ 0, 1, 2, 0, 1, 0, 0, 0 }	8238800	2687200
{ 0, 1, 1, 2, 0, 0, 0, 0 }	295680	21288960
{ 0, 1, 0, 1, 1, 0, 0, 0 }	506880	59127360
{ 0, 1, 0, 0, 1, 1, 0, 0 }	3520	5740160
{ 0, 0, 3, 1, 0, 0, 0, 0 }	14080	31680
{ 0, 0, 1, 1, 0, 1, 0, 0 }	12593688	53320
{ 0, 0, 1, 0, 1, 0, 1, 0 }	50688	12525048
	12672	202752
	-	-
{ 14, 0, 0, 0, 0, 0, 0, 0 }	32157818129473368835200	344900407873708987200
{ 12, -1, 0, 0, 0, 0, 0, 0 }	29094197009575249600	41319561145414400
{ 11, 1, 0, 0, 0, 0, 0, 0 }	-20765256335914976000	11659230882200000
{ 10, 0, 0, 0, 0, 0, 0, 0 }	235918795581162000	126823388017277200
{ 10, 1, 0, 0, 0, 0, 0, 0 }	1007050070500000	1007050070500000
{ 9, -1, 0, 0, 0, 0, 0, 0 }	1753631734111785000	74539774108000
{ 9, 1, 0, 0, 0, 0, 0, 0 }	7070599854279561600	27304174077481944000
{ 8, 1, 0, 0, 0, 0, 0, 0 }	3249373522451200	106319407748102400
{ 8, 1, 1, 0, 0, 0, 0, 0 }	544213049516163600	15142368849315200
{ 8, 0, 0, 2, 0, 0, 0, 0 }	16812561600	13532873953600
{ 8, 0, 0, 1, 0, 0, 0, 0 }	77841015997854400	2228528534897364800
{ 7, 2, 1, 0, 0, 0, 0, 0 }	3022010759400	84485590137600
{ 7, 1, 0, 0, 0, 0, 0, 0 }	3241904252145560	6439473524430320
{ 7, 0, 1, 1, 0, 0, 0, 0 }	1125538894892412000	10897284400
{ 6, 4, 0, 0, 0, 0, 0, 0 }	159775428889200	3070844243064374400
{ 6, 2, 1, 0, 0, 0, 0, 0 }	3423416081012980	6834807164303200
{ 6, 1, 2, 0, 0, 0, 0, 0 }	342301410400	32354302789800
{ 6, 1, 1, 2, 0, 0, 0, 0 }	45161977114280	6675955028160
{ 6, 0, 0, 0, 0, 0, 0, 0 }	6594496705620	85720568731520
{ 5, 6, 0, 0, 2, 0, 0, 0 }	1434629398521200	283552268768793600
{ 5, 5, 1, 0, 0, 0, 0, 0 }	193616011793200	9719732076448600
{ 5, 4, 2, 0, 0, 0, 0, 0 }	1742865846693240	13886525384000
{ 5, 3, 1, 1, 0, 0, 0, 0 }	13626800	1293175176515960
{ 5, 1, 1, 0, 0, 0, 0, 0 }	7041247439364160	3814052400
{ 5, 0, 0, 3, 0, 0, 0, 0 }	35424743954280	506214527387200
{ 5, 0, 0, 1, 0, 0, 0, 0 }	261241893920	304164201600
{ 5, 0, 0, 0, 1, 0, 0, 0 }	22771547379360	202111628771600
{ 4, 5, 1, 0, 0, 0, 0, 0 }	7806195458890400	148996991409346000
{ 4, 4, 2, 0, 0, 0, 0, 0 }	72379123952243600	341702753798400
{ 4, 3, 1, 1, 0, 0, 0, 0 }	6918350718399120	9719732076448600
{ 4, 2, 2, 0, 0, 0, 0, 0 }	11707907040800	11870907040800
{ 4, 1, 2, 1, 0, 0, 0, 0 }	2633265016640	24350653387840
{ 4, 1, 1, 0, 1, 0, 0, 0 }	227053405200	2013420727720
{ 4, 0, 1, 0, 2, 0, 0, 0 }	2335412300240	21979914259720
{ 4, 0, 0, 2, 1, 0, 0, 0 }	9777322265920	21979914259720
{ 4, 0, 0, 1, 0, 1, 0, 0 }	605558157920	157551200
{ 4, 0, 0, 0, 1, 0, 1, 0 }	853377520	45053600
{ 4, 0, 0, 0, 0, 2, 0, 0 }	1720149288	22018464468
{ 4, 0, 0, 0, 0, 1, 0, 0 }	2420714926560	603571239654100
{ 3, 1, 0, 0, 1, 0, 0, 0 }	73323572200	22791429857200
{ 3, 1, 0, 1, 0, 0, 0, 0 }	10233560216840	22383135870960
{ 3, 1, 0, 0, 0, 0, 1, 0 }	58548960	227026000
{ 3, 1, 0, 0, 0, 0, 0, 1 }	6036558157920	1224597808740
{ 3, 1, 0, 0, 0, 0, 0, 0 }	103625120	19379803680
{ 3, 1, 0, 0, 0, 0, 0, 0 }	15192777600	61769120560
{ 3, 1, 0, 0, 0, 0, 0, 0 }	747234160	434994242872
{ 3, 1, 0, 0, 0, 0, 0, 0 }	0	293434202720
{ 3, 0, 0, 0, 1, 0, 0, 0 }	21926464560	32432400
{ 3, 0, 0, 0, 1, 0, 0, 0 }	2162160	151351200
{ 3, 0, 0, 0, 0, 1, 0, 0 }	43661904	1262352400
{ 3, 0, 0, 0, 0, 0, 1, 0 }	274560	3284762980400400
{ 2, 6, 0, 0, 0, 0, 0, 0 }	55730324774200	17094673731700

(Table continued)

Table II. (*Continued*)

	50408481244261875200	10387755672522600	551317005000000000	50408481244261875200	10387755672522600	551317005000000000
9.6864476800	549174469571200	1733238888184320	3447755674522600	421817919705520	6199168621996080	50524895058159462400
142334447360	0	0	0	0	0	0
	21767607474560	613670500000000000	421817919705520	6199168621996080	50524895058159462400	50524895058159462400
1.433595566400	380540160	984951734542800	7052001816882422400	203668329324900	001144304565779200	4087712483472200
	6.3-1.0-0.0-0.0	81149691200	19439686689708480	19439686689708480	1011288779200	927504466315166770
12591419840	0	1918586689708480	0	0	0	0
	6.1-1.0-0.0-0.0	6.0-3.0-0.0-0.0	4268622883014400	4268622883014400	1011288779200	5600064683840
33902668800	0	110452835240	30719381696960	34338321362300	3297351702281355200	71346692831736200
	6.0-1.0-0.0-0.0	6.0-0.0-1.0-0.0	115452698981329900	115452698981329900	3297351702281355200	202520717143157760
-3741401664000	0	64015549973552600	30636328052000	96074955905634000	28074157094000	4843689026112100
	5.5-0.0-0.0-0.0	5.3-0-1.0-0.0	326458986603634000	326458986603634000	4843689026112100	5600064683840
-414904089600	0	1157067926868080	414904089600	96074955905634000	4843689026112100	4843689026112100
	5.2-2.0-0.0-0.0	5.1-1.0-0.0-0.0	562892805040	332224988812800	4843689026112100	4843689026112100
-36242880	0	65151604079240	65151604079240	332224988812800	4843689026112100	4843689026112100
-533210880	0	6788500000000000	6788500000000000	332224988812800	4843689026112100	4843689026112100
-1271350080	0	0	0	0	0	0
	5.0-1.0-0.0-0.0	5.0-0-1.0-0.0	57548050560	224808222480	2199207496960	2199207496960
-1361178816000	0	2210655518617200	157631040398880	300140398880	26103394437100	32599378165800
	5.0-0-0.0-0.0	5.0-0-0.0-0.0	1175869399440	1175869399440	26103394437100	26103394437100
855116890595200	0	88458362265600	3757515153733600	3872520642526200	594114247619200	574886145389045760
-65464439040	0	0	0	0	0	0
-668572183680	0	517144539341440	2077490089535360	316412599198423600	31119714949880	31119714949880
	4.1-3.0-0.0-0.0	4.1-1.0-0.0-0.0	755948793600	755948793600	3939721179920	3939721179920
345945600	0	4260683308080	10545321733440	712803817835584	65683653215360	65683653215360
	4.0-1.0-0.0-0.0	4.0-0-1.0-0.0	30320817870800	30320817870800	6182615715620	6182615715620
-1452971520	0	250403857640	637572745880	637572745880	49127572500	49127572500
	4.0-1.0-0.0-0.0	4.0-0-1.0-0.0	2421616920	7161073920	406709314371360	406709314371360
24848496	0	410329200	5805486864	87728408335974400	33581414640	33581414640
404006803200	0	7322701895045200	4353536376752600	113922686204800	87728408335974400	87728408335974400
-24014390400	0	5571518795560	763071918984905060	284037977600	3428512817114200	3428512817114200
4463513920	0	21301514795560	7185151733746	440314151301120	444638134677360	444638134677360
	3.3-2.0-0.0-0.0	3.3-1.0-0.0-0.0	62134172400	7185151733746	444638134677360	444638134677360
-129153020	0	0	0	0	0	0
	3.2-0-2.0-0.0	3.2-1-2-1.0-0.0	17315117760000	17315117760000	10335651658021120	10335651658021120
753440680	0	4754857371640	1938345621440	1020471989804080	784626168080	784626168080
	3.1-2-1.0-0.0	3.1-1-0-0.0-0.0	4237835600	1119022321280	579363697973360	579363697973360
18738720	0	6678168960	5959579509696	19224304051560	4593371476600	4593371476600
7503293720	0	75032960	7185151733746	17044302044400	19292708114200	19292708114200
4221688972	0	8958530636480	17044302044400	824933193024	449913816363200	449913816363200
	3.0-4-0-0-0-0	3.0-3-0-0-0-0	56833986198	2104161961614	12791581211712	12791581211712
11488972	0	11488972	1787963176	1079350720	144031774914	991619892400
	3.0-0-2-0-0-0	3.0-0-1-0-0-0	0	0	0	0
978737760	0	0	0	0	0	0
730527792	0	0	0	0	0	0
26565735920	0	0	0	0	0	0
	2.1-0-1-0-0-0	2.1-0-1-0-0-0	1303616156840	1437334815840	1437334815840	1437334815840
1423359396	0	0	0	0	0	0
	2.1-0-1-0-0-0	2.1-0-1-0-0-0	1423359396	13037074	1404539176400	1404539176400

(Table continued)

Table II. (*Continued*)

10

序号	项目名称	建设地点	建设性质	建设规模	主要建设内容	投资估算(万元)
1	1248640230386024000	143850024007437927880000	2389999773884601769700	3118914740880000	3118914740880000	3118914740880000
2	2064001625000000000	2064001625000000000	3631260000000000000	3631260000000000000	3631260000000000000	3631260000000000000
3	312189349797014160000	312189349797014160000	7689592745000000000	7689592745000000000	7689592745000000000	7689592745000000000
4	4388071765600000000	4388071765600000000	1625345270000000000	1625345270000000000	1625345270000000000	1625345270000000000
5	5207753305320000000	5207753305320000000	9725538398000000000	9725538398000000000	9725538398000000000	9725538398000000000
6	11087427435981760000	11087427435981760000	9309886558879424680000	9309886558879424680000	9309886558879424680000	9309886558879424680000

(Table continued)

Table II. (Continued)

{	4, 0, 0, 3, 0, 0, 0, 0,	0	1105298192000	36844114507200	385567787404900
{}	4, 0, 0, 1, 0, 0, 1, 0,	0	26854027200	19459344000	18122114000
{}	4, 0, 0, 0, 1, 0, 1, 0,	0	5139612750	66659116800	32544429120
{}	4, 0, 0, 0, 0, 2, 0, 0,	0	1857222152480000	117551796499167740	25979705482309824000
{}	3, 5, 1, 0, 0, 0, 0, 0,	0	76752285400000	439905714541000	74054938969380000
{}	3, 4, 0, 0, 0, 0, 0, 0,	0	12884498105172800	5018455800117200	811910088212422000
{}	3, 3, 1, 0, 0, 0, 0, 0,	0	1019304816200000	366725556128081200	304485641480000
{}	3, 3, 0, 0, 0, 0, 0, 0,	0	101930478076800	60304016101898000	60304016101898000
{}	3, 2, 2, 0, 0, 0, 0, 0,	0	2458174684000	7770826211724800	2098144484712600
{}	3, 2, 1, 0, 0, 0, 0, 0,	0	8986829052000	206146637720800	24832599918796800
{}	3, 2, 0, 0, 0, 0, 0, 0,	0	7042875840	299254211394080	4111519052065600
{}	3, 1, 1, 0, 0, 0, 0, 0,	0	16972413298400	36732280405687200	339066868000
{}	3, 1, 0, 1, 0, 0, 0, 0,	0	52540488000	1025512488000	668608248000
{}	3, 1, 0, 0, 1, 0, 0, 0,	0	218978760	280847827260	16961170728800
{}	3, 1, 0, 0, 0, 1, 0, 0,	0	80511883840	220622263886240	26983178490789120
{}	3, 0, 1, 0, 0, 0, 0, 0,	0	10877116421600	2061622912600	6136762956840
{}	3, 0, 1, 0, 0, 0, 0, 0,	0	4894726080	6198059147080	373325635840
{}	3, 0, 0, 1, 0, 0, 0, 0,	0	51062912560	199291521120	1273039011040
{}	3, 0, 0, 0, 1, 0, 0, 0,	0	113593345440	1519496588680	970052191350
{}	3, 0, 0, 0, 0, 1, 0, 0,	0	135933455920	5765750	86186400
{}	3, 0, 0, 0, 0, 0, 1, 0,	0	2233943513200	1457359948044000	318749408215774000
{}	3, 0, 0, 0, 0, 0, 0, 1,	0	266115514121600	74132260	2910616800
{}	2, 5, 0, 1, 0, 0, 0, 0,	0	321458443512000	1178070059168000	87310154582152000
{}	2, 4, 2, 0, 0, 0, 0, 0,	0	162458482086400	5217928063403600	5479748141868000
{}	2, 4, 1, 0, 0, 0, 0, 0,	0	15437972400	588223648000	28012362184000
{}	2, 4, 0, 1, 0, 0, 0, 0,	0	20776751151200	42353838624000	377858611138400
{}	2, 3, 3, 0, 1, 0, 0, 0,	0	98810712000	47110959589979800	4378360000
{}	2, 3, 2, 0, 0, 0, 0, 0,	0	275711436000	331862867779200	175069440300155200
{}	2, 3, 1, 0, 0, 0, 0, 0,	0	749239986705600	14992017805600	28382285131200
{}	2, 2, 2, 1, 0, 0, 0, 0,	0	2101252916000	426533572000	90480080878400
{}	2, 2, 2, 0, 1, 0, 0, 0,	0	183783600	132258646600	8414410935880
{}	2, 2, 1, 0, 1, 0, 0, 0,	0	180358647200	124410935880	3744329656010240
{}	2, 2, 0, 0, 2, 0, 0, 0,	0	24002978900	313303696149760	1194172877880
{}	2, 2, 0, 0, 1, 0, 0, 0,	0	16838355040	168737093865440	53510274446240
{}	2, 1, 1, 0, 0, 0, 0, 0,	0	1719131360	1333242330720	13938614462400
{}	2, 1, 0, 1, 0, 0, 0, 0,	0	1191361560	651910734540	13939112112950
{}	2, 1, 0, 0, 3, 0, 0, 0,	0	1045923278400	13091922043200	108004216320
{}	2, 1, 0, 0, 2, 0, 0, 0,	0	1045923278400	129129500	26885318720
{}	2, 1, 0, 0, 1, 0, 0, 0,	0	5645640	1948586860	1310400
{}	2, 1, 0, 0, 0, 1, 0, 0,	0	1498580	126117260	262080
{}	2, 1, 0, 0, 0, 0, 2, 0,	0	1498580	993982980	514920
{}	2, 1, 0, 0, 0, 0, 1, 0,	0	10822151340	5321821710600	2062260
{}	2, 1, 0, 0, 0, 0, 0, 2,	0	10822151340	5321821710600	552000
{}	2, 0, 1, 0, 0, 0, 0, 2,	0	16836920	16836920	46729866459560
{}	2, 0, 1, 0, 0, 0, 1, 0,	0	51471107200	181624400	8475667200
{}	2, 0, 1, 0, 0, 0, 0, 2,	0	51471107200	4208085800	1694364800
{}	2, 0, 1, 0, 0, 0, 0, 1,	0	5154108960	3428894480	13939112112950
{}	2, 0, 0, 1, 0, 0, 0, 1,	0	59099040	26885318720	108004216320
{}	2, 0, 0, 0, 2, 0, 0, 0,	0	59099040	4041797760	366920
{}	2, 0, 0, 0, 1, 0, 0, 0,	0	32460	514920	552000

(	1.6	1.0	0.0	0.0	0.0	0.0	)	106333547287988000	1780660844778816000
	1.5	0	0	1	0	0	)	62360511504000	7663327315584000
	1.4	1	0	1	0	0	)	9556794288934000	9556794288934000
	1.4	0	0	0	0	1	)	543358080000	543358080000
	1.3	3	0	0	0	0	)	11410780000000000	11410780000000000
	1.3	1	0	0	1	0	)	1960361288800	156719561532000
	1.3	0	0	0	0	0	)	4584918013600	472245053200
	1.2	2	0	1	0	0	)	1098779020320	1098779020320
	1.2	1	0	0	0	0	)	108013252000	1107404222400
	1.2	0	0	0	0	1	)	60510102000	42398356000
	1.2	1	0	0	1	0	)	658552000	19761424000
	1.2	0	0	1	0	1	)	16942025280	11370439800
	1.2	0	0	1	0	0	)	11643268317420	48705595360
	1.1	3	0	0	1	0	)	121827276160	106228835200
	1.1	1	0	1	0	0	)	4094226320	33636122880
	1.1	0	0	2	0	0	)	537891360	262678656240
	1.1	1	0	0	0	0	)	438197760	155568374400
	1.1	0	0	1	0	0	)	39137978880	23493166960
	1.1	0	0	0	1	0	)	5765760	114897588760
	1.1	0	0	0	0	1	)	0	80720640
	1.1	0	0	0	0	0	)	14239680	1483321280
	1.0	2	1	1	0	0	)	668107440	3423800549160
	1.0	1	3	0	0	0	)	371891520	23670015440
	1.0	1	0	1	0	0	)	7191058720	22161920
	1.0	0	1	0	1	0	)	45223735560	621505680
	1.0	0	2	0	1	0	)	263062800	491507720
	1.0	0	0	0	0	1	)	48597120	242423000
	1.0	0	0	0	0	0	)	329720	155329000
	0.9	0	0	0	0	0	)	3332875500	4813023119407500
	0.8	0	0	0	0	0	)	183725500094900	89563215085113440
	0.8	0	0	0	0	0	)	6865478000	9664778616698000
	0.7	0	0	0	0	0	)	7088482000	1523084183568000
	0.6	0	0	0	0	0	)	0	23704624944000
	0.6	0	0	0	0	0	)	16699842400	562321863200
	0.5	0	0	0	0	0	)	38160939848000	341125686864000
	0.5	0	0	0	0	0	)	0	5812593856000
	0.4	1	0	1	0	0	)	48859436000	319212437358000
	0.4	0	2	0	0	0	)	363305468000	189189000
	0.4	0	0	0	0	1	)	7068056000	90674957973600
	0.3	1	0	0	0	0	)	32612600	2547463600
	0.3	0	1	0	0	0	)	30421591200	65496772480
	0.3	0	2	0	0	0	)	82884800	1021968090
	0.3	0	0	1	0	0	)	20933732800	3977206380
	0.2	0	0	0	0	0	)	77357280	317768133920
	0.2	2	0	0	1	0	)	591714955840	55504480
	0.2	0	0	1	0	0	)	738738000	53038868800
	0.2	0	3	0	0	0	)	637831200	333091558800
	0.2	0	1	0	0	1	)	0	109108000
	0.2	0	0	0	2	0	)	3799240	108270400
	0.1	1	3	0	1	0	)	739939200	220602824
	0.1	2	0	0	0	0	)	7186222537760	2966348520
	0.1	2	0	0	0	0	)	72305513280	480480
	0.1	2	0	0	0	1	)	756755000	1021968090
	0.1	2	0	0	0	0	)	23733760	386100
	0.1	1	0	0	0	1	)	0	401849240
	0.1	0	2	0	0	0	)	25740	2720720
	0.0	0	0	0	0	0	)	0	8
)	0.0	0	0	0	0	0	)	(	)

One of us (R.Z.R.), discovered this algorithm empirically. The following beautiful proof of its correctness is due to J. Norman Bardsley, to whom the authors are very grateful.

Consider any graph with  $N$  vertices and symmetry number  $S$ . Assign, in order, the indices 1 to  $N$  to the vertices. At each stage, the next vertex can be chosen according to any topological relation that may depend on the indices previously assigned. Each ordering of the vertices  $\lambda$  is given a weight  $\omega_\lambda$ . Initially, that weight is set to unity. If at any stage in the assignment of indices the choice of the next vertex is arbitrary with degeneracy  $d$ , then  $d$  copies of the graph are made corresponding to the assignment. The weight of each copy is reduced by the factor  $d$ . When an ordering is complete, the connection matrix is constructed. Two orderings  $\lambda$  and  $\lambda'$  are equivalent if their matrices are identical. Then:

1. There are exactly  $1/S$  orderings in each equivalence class.
2. Each ordering in the same class has the same weight.
3. If  $\sum'$  represents a restricted sum over orderings including only one member of each class, then

$$\sum'_{\lambda} \omega_{\lambda} = S \quad (\text{A1})$$

since

$$1 = \sum_{\lambda} \omega_{\lambda} = \frac{1}{S} \sum'_{\lambda} \omega_{\lambda} \quad (\text{A2})$$

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